

Evolution of open quantum systems governed by unitarily covariant quantum channels

Author: Dr. Katarzyna Siudzińska

Advisor: Prof. Dariusz Chruściński

Institute of Physics, Department of Mathematical Physics
Nicolaus Copernicus University in Toruń, Poland

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Faculty of Physics, Astronomy
and Informatics

Analyzing the evolution of open quantum systems described by unitarily covariant quantum channels, including

- 1 application of the group theory to construct covariant channels,

$$U\Lambda[X]U^\dagger = \Lambda[UXU^\dagger];$$

- 2 description of quantum evolution using local and non-local master equations,

$$\dot{\Lambda}(t) = \mathcal{L}(t)\Lambda(t), \quad \dot{\Lambda}(t) = \int_0^t K(t-\tau)\Lambda(\tau) d\tau;$$

- 3 determination of Markovianity or non-Markovianity properties of quantum evolution through divisibility of dynamical maps,

$$\Lambda(t) = V(t, \tau)\Lambda(\tau), \quad \tau < t.$$

Unitarily covariant channels

Types of channels considered:

- Pauli channels (bistochastic evolution of a qubit)

$$\Lambda[X] = \sum_{\alpha=0}^3 p_{\alpha} \sigma_{\alpha} X \sigma_{\alpha}; \quad (1)$$

- Weyl channels (mixed unitary evolution)

$$\Lambda[X] = \sum_{k,l=0}^{d-1} p_{kl} W_{kl} X W_{kl}^{\dagger}; \quad (2)$$

- generalized Pauli channels

$$\Lambda[X] = p_0 X + \frac{1}{d-1} \sum_{\alpha=1}^{d+1} p_{\alpha} \sum_{k=1}^{d-1} U_{\alpha}^k X U_{\alpha}^{k\dagger}. \quad (3)$$

M. Nathanson and M. B. Ruskai, J. Phys. A: Math. Theor. **40**, 8171 (2007).

Unitarily covariant channels

A quantum channel Λ is covariant with respect to a unitary representation U of a group G if

$$\Lambda[U(g)XU^\dagger(g)] = U(g)\Lambda[X]U^\dagger(g) \quad (4)$$

for every group element $g \in G$ and Hermitian operator X .

Theorem 1

Let G be a group generated by

$$W_{kl} = \sum_{m=0}^{d-1} \omega^{km} |m + l \bmod d\rangle \langle m|, \quad \omega = e^{2\pi i/d}. \quad (5)$$

A quantum channel covariant with respect to a d -dimensional irreducible unitary representation of the group G reads

$$\Lambda[\rho] = \sum_{k,l=0}^{d-1} p_{kl} W_{kl} \rho W_{kl}^\dagger. \quad (6)$$

N. Datta, M. Fukuda, and A. S. Holevo, *Quantum Inf. Process.* **5**, 179-207 (2006).

Theorem 2

If d is a prime number, then

$$\Lambda = p_0 \mathbb{1} + \frac{1}{d-1} \sum_{\alpha=1}^{d-1} p_\alpha \mathbb{U}_\alpha \quad (7)$$

is the Weyl channel covariant with respect to **every** d -dimensional irreducible unitary representation of the group G , where

$$\mathbb{U}_\alpha[\rho] = \sum_{k=1}^{d-1} U_\alpha^k \rho U_\alpha^{k\dagger}, \quad U_\alpha = \sum_{l=0}^{d-1} \omega^l P_l^{(\alpha)}. \quad (8)$$

$P_k^{(\alpha)}$ are rank-1 projectors onto the mutually unbiased bases vectors $\{\psi_0^{(\alpha)}, \dots, \psi_{d-1}^{(\alpha)}\}$, which satisfy

$$|\langle \psi_k^{(\alpha)} | \psi_l^{(\beta)} \rangle|^2 = \begin{cases} \delta_{kl}, & \alpha = \beta, \\ \frac{1}{d}, & \alpha \neq \beta. \end{cases} \quad (9)$$

Evolution of open quantum systems

A time-evolution of a quantum system is provided by a **dynamical map** $\Lambda(t)$, $t \geq 0$, where $\Lambda(0) = \mathbb{1}$.

The simplest evolution equation for the density operator $\rho \mapsto \rho(t) = \Lambda(t)[\rho]$ is

$$\dot{\Lambda}(t) = \mathcal{L}\Lambda(t) \quad (10)$$

with the Gorini-Kossakowski-Sudarshan-Lindblad generator

$$\mathcal{L}[\rho] = -i[H, \rho] + \frac{1}{2} \sum_{\alpha} \gamma_{\alpha} \left(V_{\alpha} \rho V_{\alpha}^{\dagger} - \frac{1}{2} [V_{\alpha}^{\dagger} V_{\alpha}, \rho]_{+} \right). \quad (11)$$

Possible generalizations:

- the master equation with time-local generators $\mathcal{L}(t)$ of the form (11) but with time-dependent $H(t)$, $V_{\alpha}(t)$ and $\gamma_{\alpha}(t)$;
- the Nakajima-Zwanzig equation with a memory kernel $K(t, \tau)$,

$$\dot{\Lambda}(t) = \int_0^t K(t, \tau) \Lambda(\tau) d\tau. \quad (12)$$

Evolution of open quantum systems

Consider the generalized Pauli channel $\Lambda(t)$ generated by

$$\mathcal{L}(t) = \frac{1}{d} \sum_{\alpha=1}^{d+1} \gamma_{\alpha}(t) [\mathbb{U}_{\alpha} - (d-1)\mathbb{I}]. \quad (13)$$

Its eigenvalues are given by

$$\lambda_{\alpha}(t) = \exp[\Gamma_0(t) - \Gamma_{\alpha}(t)], \quad \Lambda(t)[U_{\alpha}^k] = \lambda_{\alpha}(t) U_{\alpha}^k, \quad (14)$$

with $\Gamma_{\alpha}(t) = \int_0^t \gamma_{\alpha}(\tau) d\tau$ and $\gamma_0(t) = \sum_{\alpha=1}^{d+1} \gamma_{\alpha}(t)$.

Theorem 3

A generator $\mathcal{L}(t)$ describes a physical evolution if and only if

$$\sum_{\alpha=1}^{d+1} e^{\Gamma_{\alpha}(t)} \leq e^{\Gamma_0(t)} + d \min_{\beta>0} e^{\Gamma_{\beta}(t)}. \quad (15)$$

Markovianity of quantum evolution

A dynamical map

$$\Lambda(t) = V(t, s)\Lambda(s) \quad (16)$$

is

- **CP-divisible** if and only if $V(t, s)$ is completely positive;
- **P-divisible** if and only if $V(t, s)$ is positive but not completely positive.

Definition 1

A **Markovian** evolution is given by a CP-divisible $\Lambda(t)$.

The dynamical map generated by the time-local generator

$$\mathcal{L}(t)[\rho] = -i[H(t), \rho] + \frac{1}{2} \sum_{\alpha} \gamma_{\alpha}(t) \left(V_{\alpha}(t)\rho V_{\alpha}^{\dagger}(t) - \frac{1}{2}[V_{\alpha}^{\dagger}(t)V_{\alpha}(t), \rho]_{+} \right) \quad (17)$$

is CP-divisible if and only if $\gamma_{\alpha}(t) \geq 0$.

D. Chruściński and A. Kossakowski, J. Phys. B: At. Mol. Opt. Phys. **45**, 154002 (2012).

Definition 2

An evolution is **weakly non-Markovian** if $\Lambda(t)$ is P-divisible.

D. Chruściński and S. Maniscalco, Phys. Rev. Lett. 112, 120404 (2014).

Theorem 4

If the Weyl channel is P-divisible, then

$$\frac{d}{dt} |\lambda_{kl}(t)| \leq 0. \quad (18)$$

Theorem 5

If the generalized Pauli channel is P-divisible, then

$$\sum_{\alpha \neq \beta} \gamma_{\alpha}(t) \geq 0. \quad (19)$$

Theorem 6

An invertible Weyl channel is P-divisible if any sum of d decoherence rates $\gamma_{kl}(t)$ is non-negative.

Theorem 7

An invertible generalized Pauli channel is P-divisible if

$$\forall_{\alpha \neq \beta} \quad \gamma_{\alpha}(t) + (d - 1)\gamma_{\beta}(t) \geq 0. \quad (20)$$

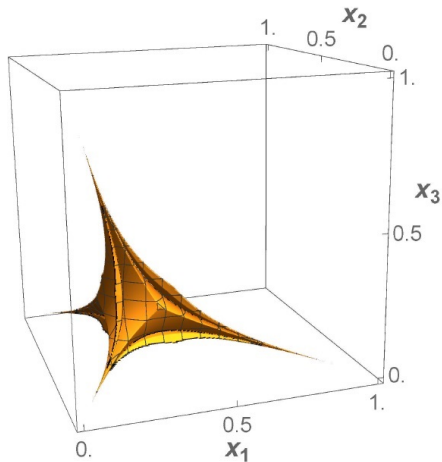
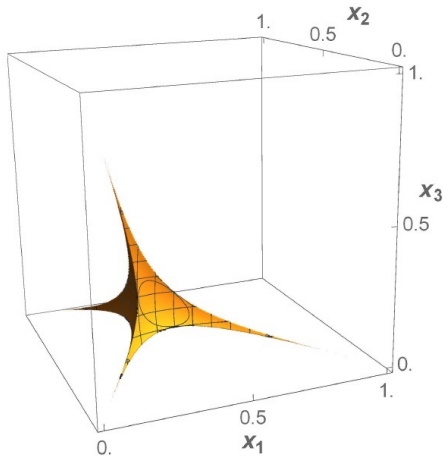
Less restrictive conditions read

$$\forall_{\gamma_{\beta}(t) \geq 0} \quad [d - 2(k - 1)]\gamma_{\beta}(t) + [d + 2(k - 1)]\gamma_{\alpha}(t) \geq 0, \quad (21)$$

where $k \leq (d + 1)/2$.

Markovianity of quantum evolution

$$\Lambda(t) = \sum_{\alpha=1}^4 x_{\alpha} e^{3t\mathcal{L}_{\alpha}} = \frac{1}{3} \left[(1 + 2e^{-3t})\mathbb{1} + (1 - e^{-3t}) \sum_{\alpha=1}^4 x_{\alpha} \mathbb{U}_{\alpha} \right]$$



Memory kernels

Consider the evolution given by the generalized Pauli channel that solves the equation

$$\dot{\Lambda}(t) = (K * \Lambda)(t) := \int_0^t K(t - \tau) \Lambda(\tau) d\tau \quad (22)$$

with the memory kernel

$$K(t) = \frac{1}{d} \sum_{\alpha=1}^{d+1} k_{\alpha}(t) [\mathbb{U}_{\alpha} - (d-1)\mathbb{I}]. \quad (23)$$

$K(t)$ has the same eigenvectors as $\Lambda(t)$,

$$\Lambda(t)[U_{\alpha}^k] = \lambda_{\alpha}(t) U_{\alpha}^k, \quad K(t)[U_{\alpha}^k] = \kappa_{\alpha}(t) U_{\alpha}^k, \quad (24)$$

and therefore the evolution equation is equivalent to

$$\dot{\lambda}_{\alpha}(t) = \int_0^t \kappa_{\alpha}(t - \tau) \lambda_{\alpha}(\tau) d\tau. \quad (25)$$

Theorem 8

A memory kernel $K(t)$ generates a physical evolution if and only if

$$\tilde{\kappa}_\alpha(s) = -\frac{s\tilde{\ell}_\alpha(s)}{1 - \tilde{\ell}_\alpha(s)}, \quad \text{where} \quad \lambda_\alpha(t) = 1 - \int_0^t \ell_\alpha(\tau) d\tau, \quad (26)$$

where the functions $\ell_\alpha(t)$ satisfy the conditions

$$\int_0^t \ell_\alpha(\tau) d\tau \geq 0, \quad (27)$$

$$d \int_0^t \ell_\beta(\tau) d\tau \leq \sum_{\alpha=1}^{d+1} \int_0^t \ell_\alpha(\tau) d\tau \leq \frac{d^2}{d-1}. \quad (28)$$

The Laplace transform of $f(t)$: $\tilde{f}(s) := \int_0^\infty f(t)e^{-st} dt$.

J. Williams, Laplace Transforms, George Allen & Unwin, London 1973.

Definition 3

Two families of completely positive maps $\{N(t), Q(t)\}$ form a **legitimate pair** if

$$N(0) = \mathbb{1}, \quad \|\tilde{Q}(s)\|_1 < 1, \quad \text{Tr } Q(t)[X] + \text{Tr } \dot{N}(t)[X] = 0. \quad (29)$$

A legitimate pair allows one to construct the dynamical map

$$\Lambda(t) = N(t) + (N * Q)(t) + (N * Q * Q)(t) + \dots, \quad (30)$$

which is the solution if the master equation with the memory kernel

$$\tilde{K}(s) = s\mathbb{1} - \left[1 - \tilde{Q}(s)\right] \tilde{N}^{-1}(s). \quad (31)$$

D. Chruściński and A. Kossakowski, Phys. Rev. A **94**, 020103(R) (2016).

Example – beyond the legitimate pairs

The generalized Pauli channel for which $Q(t)$ is **not** completely positive:

$$\Lambda(t) = \frac{1}{d+1-k} \sum_{\alpha \notin \mathcal{J}} e^{dt\mathcal{L}_\alpha}, \quad \mathcal{J} \subset \{1, \dots, d+1\}, \quad |\mathcal{J}| = k. \quad (32)$$

Memory kernels

A **semi-Markov evolution** is fully determined by a family of completely positive maps $Q(t)$ for which

$$\int_0^t Q^\dagger(\tau)[\mathbb{I}] d\tau \leq \mathbb{I}, \quad \text{Tr}(XQ(t)[Y]) =: \text{Tr}(Q^\dagger(t)[X]Y). \quad (33)$$

For the generalized Pauli channels, one finds

$$Q(t) = \frac{1}{d-1} \sum_{\alpha=1}^{d+1} f_\alpha(t) \mathbb{U}_\alpha, \quad N(t) = \left[1 - \sum_{\alpha=1}^{d+1} \int_0^t f_\alpha(\tau) d\tau \right] \mathbb{I}. \quad (34)$$

D. Chruściński and A. Kossakowski, Phys. Rev. A **95**, 042131 (2017).

Example

A semi-Markov evolution is provided by the convex combination of Markovian semigroups

$$\Lambda(t) = \frac{1}{d+1} \sum_{\alpha=1}^{d+1} e^{dt\mathcal{L}_\alpha}, \quad f_\alpha(t) = \frac{d-1}{d+1} \exp\left[-\frac{d(d+1)-1}{d+1}t\right] \quad (35)$$

The most important results:

- construction of quantum channels by imposing the condition of covariance with respect to representations of finite groups;
- necessary and sufficient conditions for admissible generators and memory kernels;
- P-divisibility conditions for the generalized Pauli channels.

Further work:

- mixtures of Markovian generalized Pauli dynamical maps;
- geometry of Pauli maps and generalized Pauli channels;
- generalization of Pauli channels using mutually unbiased measurements.

Ending remarks

This PhD thesis was based on the following publications:

- 1 D. Chruściński and K. Siudzińska, *Generalized Pauli channels and a class of non-Markovian quantum evolution*, Phys. Rev. A **94**, 022118 (2016);
- 2 K. Siudzińska and D. Chruściński, *Memory kernel approach to generalized Pauli channels: Markovian, semi-Markov, and beyond*, Phys. Rev. A **96**, 022129 (2017);
- 3 K. Siudzińska and D. Chruściński, *Quantum channels irreducibly covariant with respect to the finite group generated by the Weyl operators*, J. Math. Phys. **59**, 033508 (2018).

The PhD thesis was created during the realization of research projects

- 1 *Non-Markovian quantum dynamics* No. 2015/17/B/ST2/02026 (OPUS 9),
- 2 *Evolution of open quantum systems governed by unitarily covariant quantum channels* No. 2018/28/T/ST2/00008 (ETIUDA 6)

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